

# Short-Term Tidal Variability During Sudden Stratospheric Warming

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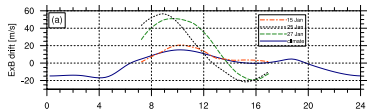
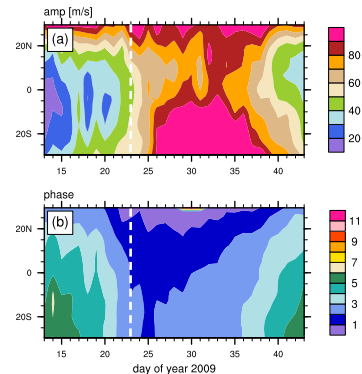
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# SW2 amplitude and phase changes during SSW 2009

(Wang et al., 2014)



Amplitude of zonal wind @ 115 km

+

Phase of zonal wind @ 115 km

↓

Plasma drift @ Jicamarca from GIP  
GIP=Global Ionosphere Plasmasphere

What can cause SW2 amplitude and phase changes during SSW 2009?

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# What can cause the short-term variability of tides?

The short-term variability of tides can be caused by (e.g., Vial 1993)

- Change in tidal forcing;
- Change in background propagation conditions;
- Introduction of energy near tidal frequencies by local or synoptic scale disturbances; and
- Non-linear interactions with atmospheric waves of different scales (and frequency, including modulation by planetary waves).

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# What can cause tidal changes during SSW?

Several mechanisms have been proposed:

- ① Stratospheric ozone change  $\Rightarrow$  heating changes (e.g., Goncharenko et al., 2012);
- ② Reduction of north-south asymmetry of zonal background wind (Jin et al., 2012);
- ③ Nonlinear interactions between tides, and/or with planetary waves (e.g., Wang et al., 2012; Liu et al., 2010); and
- ④ Change in lunar tides (e.g., Fejer et al., 2010; Forbes and Zhang, 2012; Pedatella et al., 2014).

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## Objective and methods

- To understand the short-term variability of tides during SSW; in particular the cause of SW2 amplitude and phase changes.
- To achieve this goal, we carry out a detailed analysis of WAM simulation of SSW 2009. In particular, we
  - Consider the first-order effects of the background wind and temperature;
  - Decompose the perturbation fields into Fourier modes in longitude-time and Hough modes in latitude; and
  - Solve the tidal vertical structure equation (VSE).

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## The equations for perturbations

$$\text{u-momentum} \quad \frac{\partial u'}{\partial t} - fv' + \frac{1}{a \cos \phi} \frac{\partial \Phi'}{\partial \lambda} = R'_u, \quad (1a)$$

$$\text{v-momentum} \quad \frac{\partial v'}{\partial t} + fu' + \frac{1}{a} \frac{\partial \Phi'}{\partial \phi} = R'_v, \quad (1b)$$

$$\text{thermodynamic} \quad \frac{\partial \theta'}{\partial t} + \frac{\partial \bar{\theta}}{\partial z} w' = Q' + R'_\theta, \quad (1c)$$

where

$$R'_u \equiv -\frac{\bar{u}}{a \cos \phi} \frac{\partial u'}{\partial \lambda} - \frac{v'}{a \cos \phi} \frac{\partial \bar{u} \cos \phi}{\partial \phi} - \frac{\partial \bar{u}}{\partial z} w', \quad (2a)$$

$$R'_v \equiv -\frac{\bar{u}}{a \cos \phi} \frac{\partial v'}{\partial \lambda} - \frac{2\bar{u} \tan \phi}{a} u', \quad (2b)$$

$$R'_\theta \equiv -\frac{\bar{u}}{a \cos \phi} \frac{\partial \theta'}{\partial \lambda} - \frac{1}{a} \frac{\partial \bar{\theta}}{\partial \phi} v'. \quad (2c)$$

## Fourier expansion in longitude and time

$$u'(\lambda, \phi, z, t) = \tilde{u}(\phi, z) \exp[i(2\Omega\sigma t + s\lambda)], \quad (3)$$

$$\begin{aligned} \frac{\partial^2 \tilde{w}}{\partial z^2} - \frac{1}{H} \frac{\partial \tilde{w}}{\partial z} - \frac{N^2}{4a^2\Omega^2} \mathcal{F}(\tilde{w}) = & -\frac{1}{4a^2\Omega^2} \frac{R}{H} e^{-\kappa z/H} \mathcal{F}(\tilde{Q} + \tilde{R}_\theta) \\ & - \frac{\sigma}{2sa\Omega} \mathcal{F}_u \left( \frac{\partial \tilde{R}_u}{\partial z} \right) + \frac{i\sigma}{2a\Omega} \mathcal{F}_v \left( \frac{\partial \tilde{R}_v}{\partial z} \right), \end{aligned}$$

$$\begin{aligned} \mathcal{F}(A) \equiv & \frac{1}{\cos \phi} \frac{\partial}{\partial \phi} \left( \frac{\cos \phi}{\sigma^2 - \sin^2 \phi} \frac{\partial A}{\partial \phi} \right) \\ & - \frac{1}{\sigma^2 - \sin^2 \phi} \left( \frac{s \sigma^2 + \sin^2 \phi}{\sigma \sigma^2 - \sin^2 \phi} - \frac{s^2}{\cos^2 \phi} A \right), \end{aligned} \quad (4a)$$

$$\mathcal{F}_u(A) \equiv \frac{1}{\cos \phi} \left[ \frac{\partial}{\partial \phi} \left( \frac{s \cos \phi \sin \phi}{\sigma \sigma^2 - \sin^2 \phi} A \right) + \frac{s^2}{\sigma^2 - \sin^2 \phi} A \right], \quad (4b)$$

$$\mathcal{F}_v(A) \equiv \frac{1}{\cos \phi} \left[ \frac{\partial}{\partial \phi} \left( \frac{\cos \phi}{\sigma^2 - \sin^2 \phi} A \right) + \frac{s}{\sigma} \frac{\sin \phi}{\sigma^2 - \sin^2 \phi} A \right]. \quad (4c)$$

# Expansion in Hough modes

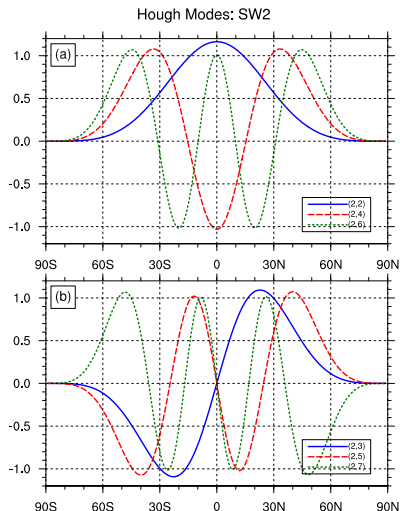
- Hough modes are eigenfunctions of Laplace tidal operator

$$\mathcal{F}(\Theta_n) + \gamma_n \Theta_n = 0.$$

- Expansion in Hough modes

$$\hat{w}_n(z) = \int_{-1}^1 \hat{w}(\mu, z) \Theta_n(\mu) d\mu,$$

where  $\mu = \sin \phi$  and the integration is done via Gauss quadrature rule.



## The tidal VSE

We now have the tidal VSE in its final form (an ODE!)

$$\frac{d^2 \hat{w}_n}{dz^2} + \left( \frac{\gamma_n \bar{N}^2}{4a^2 \Omega^2} - \frac{1}{4H^2} \right) \hat{w}_n = \frac{\gamma_n}{4a^2 \Omega^2} \frac{R}{H} e^{-\kappa z/H} \hat{Q}_n - \frac{\sigma}{2sa\Omega} \hat{F}_n^u + \frac{i\sigma}{2a\Omega} \hat{F}_n^v + \frac{\gamma_n}{4a^2 \Omega^2} \frac{R}{H} e^{-\kappa z/H} \hat{R}_{\theta n}, \quad (5)$$

where

$$\hat{F}_n^u \equiv \int_{-1}^1 \mathcal{F}_u \left( \frac{\partial \hat{R}_u}{\partial z} + \frac{\hat{R}_u}{2H} \right) \Theta_n(\mu) d\mu, \quad (6a)$$

$$\hat{F}_n^v \equiv \int_{-1}^1 \mathcal{F}_v \left( \frac{\partial \hat{R}_v}{\partial z} + \frac{\hat{R}_v}{2H} \right) \Theta_n(\mu) d\mu, \quad (6b)$$

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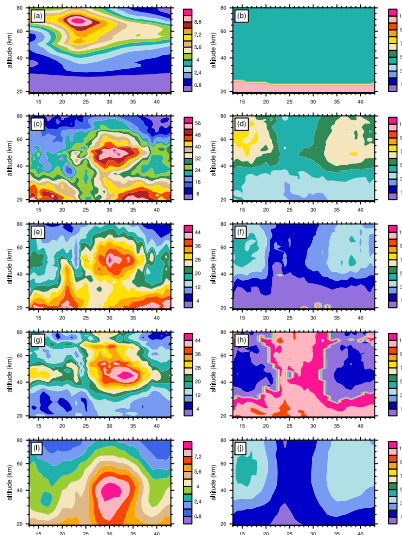
**Hough analysis**

Solve the tidal VSE

## ④ How about phase change?

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# Hough analysis for (2,2) mode



$\hat{u}_n$ : symmetry of background wind

$\hat{u}_n$ : large change in amp. & phase

$\hat{v}_n$ : 90° phase shift from  $\hat{u}_n$

$\hat{\theta}_n$ : large change in amp. & phase

$\hat{w}_n$ : large change in amp. & phase

Large consistent change in amp. & phase are shown in  $\hat{u}_n$ ,  $\hat{v}_n$ ,  $\hat{w}_n$ , &  $\hat{\theta}_n$ .



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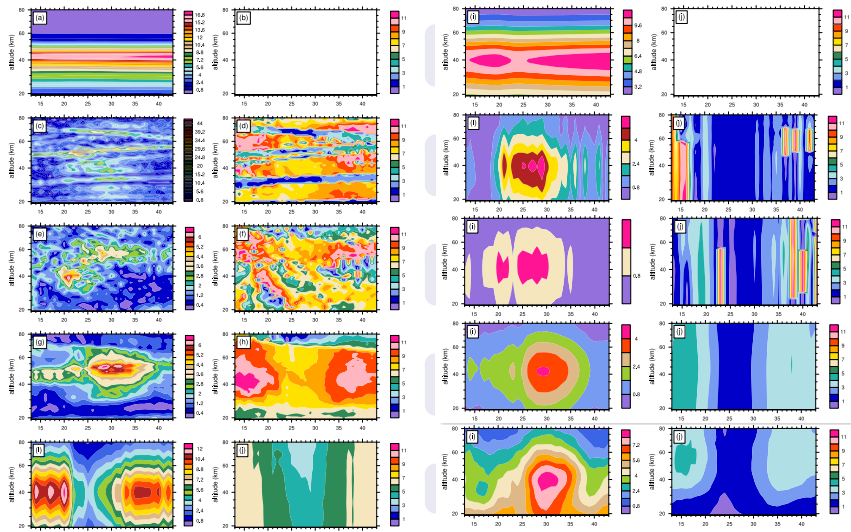
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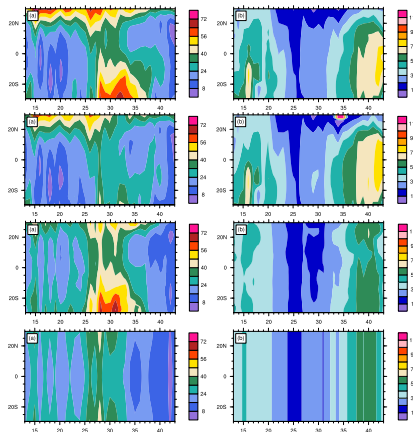
## ⑤ Summary and Conclusions

# Solve the tidal VSE for (2,2) mode



Can the linear theory explain the full nonlinear model simulation?

## Another way to solve the linear equations



All modes *with* background terms

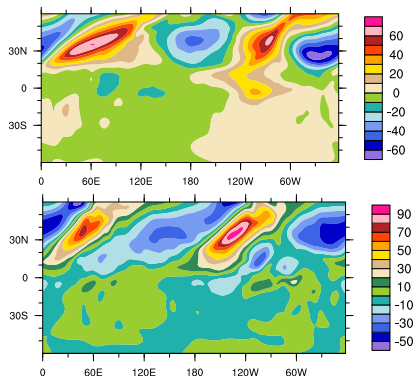
All modes *without* background

(2,2) mode *with* background

(2,2) mode *without* background

- The background terms contribute about 1/3 of the changes;
- The first symmetric (2,2) mode dominates.

# Heuristic understanding of phase changes during SSW



$u'$  at 2009011700 UTC

$u'$  at 2009012300 UTC

- The change in the prevailing stratospheric background winds during the SSW can cause also the tidal wave phase change (Wang et al. 2014);
- It can be understood as the mass or wind anomaly being pushed/sheared more westward (or early local time) *materially*.

## Summary and Conclusions

To understand causes of significant amplitude & phase changes in SW2 in WAM simulation of SSW 2009:

- A Hough function analysis shows large consistent amplitude & phase changes in SW2, mainly the first symmetric mode (2,2);
- The solution of the tidal vertical structure equation (VSE) with background wind and temperature shows:
  - The reduction of north-south asymmetry during SSW alone is not enough to explain the large amplitude & phase changes *after* the peak warming;
- The linear theory shed some lights on some features of tidal amplitude and phase change; but is not enough;
- Phase changes can be heuristically understood as the reversal of prevailing background wind “push” /shear the mass/wind anomaly more westward (or earlier local time) *materially*.